THERMAL INFRARED CONSTRAINT TO A PLANETARY COMPANION OF VEGA WITH THE MMT ADAPTIVE OPTICS SYSTEM¹

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ABSTRACT

Vega may have a massive companion in a wide orbit, as evidenced by structure in its cold dust debris. We have tested this hypothesis by direct imaging with adaptive optics in the *M* band. The observations were made with a newly commissioned thermal infrared camera, Clio, on the 6.5 MMT AO system with low-background deformable secondary mirror. The observations constrain a planet to be less than 7 M_J at the approximate position angle expected from the dust structure and at a radius >20 AU (2".5). This result is more stringent than similar previous near-infrared observations of Vega, which achieved limits of 20 and $10 M_J$ at separations of 7". The higher sensitivity is due to both the more favorable contrast of gas giant planets at the *M* band and the higher Strehl ratio and more stable point spread function at longer wavelengths. Future *L'* or *M* band observations could provide a powerful approach for wide-separation planet detection, especially for cooler and thus older or less massive planets. The natural best targets are nearby stars, where planets in the range of $5-15 M_J$ and as old as several Gyr are expected to be detectable with this technique. *Subject headings:* infrared: stars — instrumentation: adaptive optics — planetary systems — stars: individual (Vega)

1. INTRODUCTION

Direct detection of extrasolar planets is a highly desirable goal for a range of reasons. Although detection of planets through gravitational influence on the star has yielded important information about planets around other stars, many parameters of a planet are most easily derived from its spectral energy distribution (SED), including its temperature, size, and composition. For planets in orbit beyond approximately 5-10 AU, the length of time needed to detect a planet through indirect means also becomes prohibitive. Yet these are precisely the radii where we see massive planets in our own solar system.

Initial success in the area of direct detection is just now occurring, although the detectable objects are still significantly different from what we think of as a typical solar system. Measurements of secondary eclipses of transiting hot Jupiters can begin to constrain the temperature and albedo of these objects (Deming et al. 2005; Charbonneau et al. 2005; Rowe et al. 2005). For younger stars and wider separations faint companions have been detected (Chauvin et al. 2004; Neuhauser et al. 2005), which are consistent with being planetary mass objects. However, these objects are unlikely to have formed by either core accretion or gravitational instability (Boss 2006), suggesting that we are seeing either scattered planets or low-mass objects that have formed by fragmentation.

Current direct detection searches typically focus on very young stars where giant planets are early in the process of cooling and contracting and thus are expected to be relatively bright in the near-infrared (see Metchev et al. 2004; Oppenheimer et al. 2004; Mugrauer et al. 2005; Lowrance et al. 2005; and Biller et al. 2005 as examples of ongoing surveys). Of these bands, the *H* band is particularly attractive due to the methane absorption feature, which creates an identifiable spectral feature, as well as its relative brightness compared to the *K* band. The majority of the direct-detection plans with large telescopes focus in this spectral region, utilizing

techniques such as simultaneous spectral difference imaging (Marois et al. 2000; Biller et al. 2005) and higher order adaptive optics (Oppenheimer et al. 2004; Macintosh et al. 2003b) to optimize the achievable contrast. Several observatories (see Gratton et al. 2004 as an example) are currently designing ambitious nearinfrared instruments in order to pursue the detection of young planets.

Although near-infrared sensitivity from the ground is better than at longer wavelengths, a significant contrast advantage can be achieved by looking for cool giant planets closer in wavelength to the peak of their expected SED. Giant planets have very nonblackbody SEDs, requiring a knowledge of their spectral characteristics to best choose an optimum wavelength. In addition, the sensitivity of a ground-based system is dominated by the spectral dependence of the atmospheric transparency and brightness, limiting possible detection to discrete spectral windows.

It has long been known that Jupiter's SED displays an anomalous peak at \sim 4–5 μ m (Gillett et al. 1969). This is due to the lack of absorption features in this spectral region, which allows the observation of thermal emission from much deeper in the planet's atmosphere than at other wavelengths. A similar, broader peak can be seen in T-type brown dwarfs (Oppenheimer et al. 1998) and is an expected generic feature of objects with effective temperatures below approximately 1200 K, as modeled by Burrows et al. (1997, 2003) and Baraffe et al. (2003). The high flux appears in the L' and M bands for hotter objects and narrows to a feature that is well matched to the M-band window for the coolest objects. The relative flux in the L' versus M bands for hotter objects, as measured by Golimowski et al. (2004), appears to be dependent on the amount of carbon monoxide in the atmosphere due to nonequilibrium effects, lessening the expected brightness at M compared to equilibrium models. Such effects may determine whether the M or L' atmospheric window is preferable, but the existence of a broad hump in the SED in the $4-5 \mu m$ region appears to be secure, from both model predictions and observed cool objects.

In order to make use of this expected contrast enhancement, it is necessary to be able to detect the relatively faint planet flux in

¹ Observations reported here were obtained at the MMT Observatory, a joint facility of the University of Arizona and the Smithsonian Institution.

the presence of background from the sky as well as from any warm optics in the system. The relatively small expected separations require diffraction-limited performance of the optics, which in turn requires an adaptive optics system to correct for atmospheric turbulence. Typical AO systems add 5–10 warm surfaces. For wavelengths longward of $\sim 2 \mu m$ the infrared glow from the warm optics can dominate the background and thus the noise of an optical system. For planet detection at longer wavelengths, an optimum system is one that is integrated into the telescope itself, such as the deformable secondary mirror of the MMT Adaptive Optics (AO) system (Riccardi et al. 2002; Brusa et al. 2003).

Optimal targets for direct detection of extrasolar planets are stars that are both young and nearby. Unfortunately, most radial velocity targets do not satisfy both of these criteria. Target stars are typically chosen to avoid potential noise in a Doppler signal from stellar activity (Wright et al. 2004). For young stars, one possible alternative indicator of the existence of a planetary companion is substructure in a debris disk (Roques et al. 1994; Liou & Zook 1999). Giant planets will tend to set up resonances in debris material around a star, with the characteristics of the resonances dependent on parameters of the planets' mass and orbit (Kuchner & Holman 2003). If this substructure can be modeled properly, it may be possible to use such information to guide where and which stars to search for planetary companions.

In this paper we present initial observations with Clio (Freed et al. 2004), a camera designed specifically to detect giant planets through their $3-5 \ \mu m$ infrared radiation. We have used the imager to constrain the existence of a planetary companion to Vega at the orientation expected from resonances in the cold debris disk. Section 2 describes the instrumental setup. In § 3 we present the observations. Section 4 details the data reduction and achievable limits versus separation. In § 5 we discuss the implications of these initial observations for future planet detection.

1.1. The Vega System

Vega provides an interesting target for planet searches, as one of the nearest young stars with a debris disk indicative of planetesimals. For the purpose of planet models used in this paper we adopt an age of 300 Myr and a distance of 7.8 pc, consistent with evolutionary tracks (Song et al. 2001) and *Hipparcos* astrometry. The debris disk around Vega has been intensively studied since its discovery with IRAS (Aumann et al. 1984). It appears that we are viewing a cold (Backman & Paresce 1993), probably transiently bright (Su et al. 2005) dust disk equivalent to the Kuiper Belt at nearly pole-on orientation around Vega. Submillimeter and millimeter observations of Vega have revealed a double-lobed enhancement to the ring (Holland et al. 1998; Koerner et al. 2001; Wilner et al. 2002, hereafter W02). The offset nature of the lobes seen at long wavelengths has been modeled by W02. They reproduce the observed structure for a planet that is $\sim 90^{\circ}$ from the line connecting the lobes on the same side of the star as the offset of the lobes. They suggest that the planet is 7" from the star in the northwest direction.

Interestingly, the observed structure seen around Vega is similar to one of the four broad classes of dust structure described by Kuchner & Holman (2003), which can arise from giant planet perturbations to debris disks. The class that fits the double-lobed pattern of Vega is indicative of a high-mass planet on an eccentric orbit. The mass of the planet is not well constrained by the observable resonant structure, although W02 achieve good correspondence between their model and the observations using a $3 M_J$ planet. Another hypothesis for the lobed structure has been proposed by Wyatt (2003), where the resonance arises from outward

planetary migration of a less massive (Neptune-mass) planet from 40 to 65 AU. Direct imaging has the potential to detect massive planets around Vega, potentially providing an important constraint to modeling of the resonant structure.

Detection of a companion to Vega has previously been attempted with the Palomar system in the *H* band. Metchev et al. (2003) set a 5 σ limit of approximately 20 M_J based on their observations at the separation of the expected perturber. Macintosh et al. (2003a) also used Keck at the *K* band to carry out a similar observation. They estimate their 5 σ limit corresponded to 10 M_J at the expected separation of the planet. More recent deep integrations have been achieved by Marois et al. (2006), reaching a limit of ~4 M_J at a separation of 7".

2. INSTRUMENTAL SETUP

The observations described below made use of two recent developments for the MMT telescope: a deformable secondary adaptive optics system and a high well depth, high duty cycle indium antimonide based camera system. We describe the advantages of these approaches below.

2.1. The MMT Deformable Secondary AO System

The MMTAO system makes use of the world's first deformable secondary mirror (Riccardi et al. 2002; Brusa et al. 2003), allowing integration of the AO into the telescope itself (Wildi et al. 2003). The deformable secondary is undersized to form the stop of the system, and provides an effective aperture of 6.35 m diameter. Secondary undersizing is a proven technique to provide efficient baffling of warm background radiation from the telescope structure. Infrared light encounters only two warm reflecting surfaces, the primary and secondary mirror, before entering the cryogenically cooled camera, which eliminates background contributions from these optics. The result is an infrared optimized system that is also capable of diffraction-limited imaging (a full width at half-maximum of 0.13 at L' and 0.16 at the *M* band). Compared to conventional adaptive optics systems, this results in improved sensitivity. Lloyd-Hart (2000) shows that for typical parameters expected at the MMT, observations with an optimized system can achieve the same signal in onethird to one-half of the integration time that would otherwise be needed using a conventional adaptive optics system at L' and M bands.

Atmospheric turbulence is sensed using visible light reflected from the entrance window of the cryostat. The telescope pupil is imaged onto a CCD-based Shack Hartmann wave-front sensor operating at 550 Hz. The resulting slopes are reconstructed to measure and correct the first 56 Zernike modes of the turbulent wave front. The system typically delivers Strehl ratios of $\sim 25\%$ at the *H* band in median seeing (Miller et al. 2004). Scaled to *M* band, this is a Strehl ratio of 85%.

2.2. Clio: A 3-5 µm Imager

Clio (Freed et al. 2004; Sivanandam et al. 2006) is an imager designed for obtaining high spatial resolution images with optimum efficiency at the L' and M bands. The favorable contrast for planets at the M band is at least partially offset by the high background in this atmospheric window. Our initial estimates of the relative sensitivity in these two bands suggests that the L'band magnitude limit is ~2 mag better than the M-band limit, consistent with the higher background at the M band compared to L'. This is a good match to the expected colors for giant planets. For example, models of giant planets (Baraffe et al. 2003; Burrows et al. 2003) predict that a planet will be roughly 2 mag brighter in the M band than in L' for an effective temperature around 350 K. The coincidence of these two factors, a fainter L'limiting magnitude, and a red L'-M color for a typical planet, provides favorable conditions for confirmation of any substellar companion. While common proper motion will still be important to verify physical association, confirmation via color is a strong secondary indicator that an object is indeed a planetary mass object.

High-efficiency observations with Clio are aided by an indium antimonide (InSb) detector from Indigo Systems, Inc., which has a high well depth compared to more typical astronomical InSb detectors. This feature allows efficient use in the *M* band, where high backgrounds can typically swamp detectors with lower well depths. The detector has a measured well depth of 3 million photoelectrons (Sivanandam et al. 2006). For the initial observations, the frame rate was only 4 Hz. This resulted in a duty cycle of roughly 50% for the 120 ms frame time used in the observations of Vega. The camera readout has been since improved to 20 Hz. This will allow a higher observational efficiency with Clio in the future.

The camera optics reimage the telescope to a cold mask, which eliminates the surrounding radiation from warm telescope structure. The f/15 Cassegrain focus is reimaged at f/20 onto the 320 × 256 pixels, each 30 μ m in width. The resulting plate scale is ~0".05 per pixel. The detector is oriented so that the long direction is in the elevation direction when the rotator offset is zero.

3. OBSERVATIONS

Observations of Vega were taken with Clio and the MMT AO system on 2005 June 21, during initial tests of the Clio camera. The filter used was from OCLI (now JDS Uniphase) with halfpower points at 4.35 and 4.95 μ m, slightly shorter than a standard *M*-band filter. L' band images were also taken, but were later found to be contaminated by a filter misalignment in the camera and are not presented here. Table 1 lists the data taken. Individual images, each 120 ms integration time, were co-added in computer memory. One hundred images were combined for each frame, resulting in an integration time of 12 s for each frame. Fifty-six frames were acquired in closed loop with sufficiently high Strehl ratio to use in the final image, resulting in a total integration time of 672 s. The telescope was periodically nodded by 5" in the direction to move Vega vertically on the array to allow for background subtraction. The star was placed on the lower left or upper left quadrant of the array for the two positions. This allowed roughly 220 pixels or 11" of field on the west side of the star, the location proposed for the hypothetical perturbing companion. The instrument derotator was turned off, which resulted in slow sky rotation of the field. Since we observed Vega well after transit, this amounted to no more than 10° rotation during the course of all the observations. For roughly half the frames the derotator was turned from an offset of zero to an offset of -30° . This was done to check the approach of PSF subtraction via rotator offset.

4. DATA REDUCTION AND ANALYSIS

Data analysis was carried out with custom C software developed for reduction of Clio data. The software is being developed to provide relatively automated reduction of Clio data optimized for faint companion detection. Initial background removal was achieved via subtraction of each frame with its nod pair. Although the images were saturated inward of approximately 0."3, boxcar smoothing the resulting images with a 25×25 pixel aperture allowed a valid centroid to be calculated. The calculated centroid for the positive and negative nod image in each frame was used to shift each image. The parallactic angle was calcu-

 TABLE 1

 Log of Data Acquisition for Clio Observations of Vega 2005 June 21

Frames	Integration Time (s)	Position	Parallactic Angle	Rotator Offset (deg)	Comment
Traines	(3)	1 0311011	Aligic	(ueg)	Comment
1-10	12	upper left	-267	0	
11-20	12	lower left	-268	0	No AO last two
21-22	12	lower left	-270	0	
23–24	12	upper left	-270	0	
25-26	12	lower left	-271	0	
27-28	12	upper left	-271	0	
29-30	12	lower left	-271	0	
31-32	12	upper left	-272	0	
33-34	12	lower left	-272	0	
35-36	12	upper left	-273	0	
37-38	12	lower left	-273	0	
39-40	12	upper left	-273	0	
41-42	12	lower left	-273	0	
43-44	12	upper left	-273	0	
45-46	12	lower left	-273	0	
47-48	12	upper left	-273	0	
49-50	12	lower left	-273	0	
51-52	12	upper left	-276	-30	
53-54	12	lower left	-276	-30	
55-56	12	upper left	-276	-30	
57-58	12	lower left	-276	-30	No AO
59-60	12	lower left	-276	-30	No AO
61–62	12	lower left	-276	-30	
63–64	12	upper left	-276	-30	
65-66	12	lower left	-276	-30	
67–68	12	upper left	-276	-30	
69–70	12	lower left	-276	-30	

lated from the time taken for each frame and used to derotate the images before final combination. Additional processing to account for drifts in detector bias and noise in individual columns was carried out and is described further by Heinze et al. (2006). The resulting image has a total integration time on Vega of 672 s. A broad halo and diffraction pattern from Vega dominates the resulting image. To remove this, the image was boxcar smoothed using a Gaussian kernel with FWHM of 5 pixels, and the resulting image. Figure 1 shows the resulting image. The most striking feature of the data is the number of Airy rings detectable, a characteristic of the high Strehl ratio achieved at the *M* band with adaptive optics.

Noise contours were generated from the image in order to quantify where we are sensitive to what level of companions; these are shown in Figure 2. The boundaries between each successive gray scale are, from faintest to brightest, 1.13, 1.72, 2.54, 16, and 100 mJy. For the models of Burrows et al. (2003), this corresponds to a limit of 7, 10, 14, and $26 M_J$ for the four lowest contours, respectively. For regions of overlap of all of the frames the \sim 5 σ limit to a planet is 0.78 mJy, which corresponds to a mass limit of 6 $M_{\rm J}$. The data show that the noise is dominated by sky background outside of approximately 2".5. Figure 1 also shows the insertion of synthetic planets into the data. Inspection of these planet images reveals that the actual limiting magnitude might be slightly higher than the 0.8 mJy formal calculation, possibly due to some amount of correlation in the detector pixels, an artifact of column noise that is only partially removable. However, the 1.1 mJy sources appear detectable, allowing us to feel confident that we would detect a planet down to approximately 7 $M_{\rm I}$ in the darkest contour shown in Figure 2.

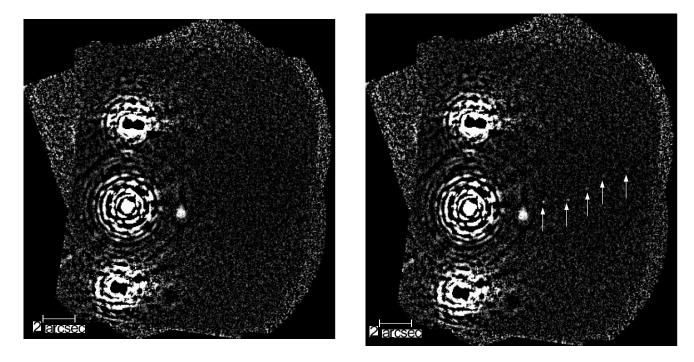


Fig. 1.—Co-added image of Vega. The left panel shows a co-added image with a total integration time of 672 s. The image has been unsharp masked, which reveals the Airy pattern diffraction as well as a ghost image to the right of the star. The artifacts approximately 5'' above and below the star are residuals of the nod subtraction. North is up and east to the left. No planets are seen to the limit of our sensitivity. The planet hypothesized by W02 would be in the upper right quadrant of this image. The right panel is the same image with five synthetic planets placed in the raw data, so that they appear at the position of the arrows. From left to right, the planets are the brightness expected for a 10, 7, 7, 7, and 6 $M_{\rm I}$ planet. The rightmost planet is formally at our sensitivity limit, although it cannot be seen reliably in the image.

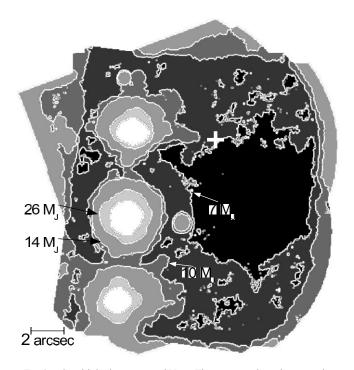


FIG. 2.—Sensitivity image around Vega. The contours show the approximate sensitivity limit for different regions of the image. From brightest to faintest, the contours correspond to sensitivity limits (and corresponding mass limits) of 100 mJy, 16 mJy ($26 M_J$), 2.54 mJy ($14 M_J$), 1.72 mJy ($10 M_J$), and 1.13 mJy ($7 M_J$). A cross is drawn at 7" northwest of Vega, the position of a planet predicted by W02.

4.1. Optimal PSF Subtraction

To test various ways of removing the diffracted or aberrated light from the star, we investigated unsharp masking, PSF subtraction via rotator offset changes, and PSF subtraction via sky rotation. This data reduction was carried out independently of the data reduction described in \S 4, allowing an independent check on the sensitivity limit. The scripts were developed using the Perl data language.² Separate co-added frames were generated using frames 25-48 and 51-70, respectively, for images at 0 and -30° of the telescope rotator offset. No rotation of the images was carried out for this data reduction, since the main goal was to determine the stability of the PSF, not use the images for faint companion detection. Figure 3 shows the two resulting images at each rotator offset. The images were unsharp masked using a 0."25 boxcar. The static speckle pattern in the image is clearly tied to the rotator offset, suggesting the source is the telescope optics rather than the camera optics. This results in poor subtraction if the rotator offset is used, and suggests that sky rotation will be the preferred way to obtain multiple images for later PSF subtraction. This corroborates the results of Marois et al. (2006) that angular differential imaging (ADI) is a useful approach to obtain good PSF subtraction with an adaptive optics system.

Although our images were not taken during significant sky rotation, it is possible to simulate this process to evaluate the achievable limits. To do so, we rotated the second image in Figure 3 by 30° so that the static pattern of the two images are aligned. These images were taken roughly 40 minutes apart and across a change in elevation of the telescope of 7° , parameters similar to what would be expected when taking multiple images during a rapid change in parallactic angle of the object. Good subtraction of the PSF is obtained using this approach.

² See http://www.pdl.perl.org.

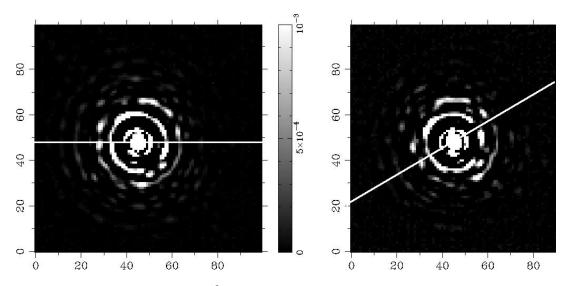
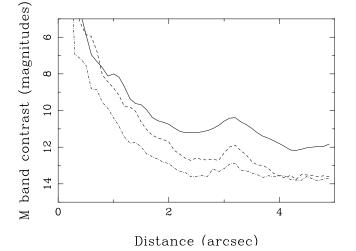


Fig. 3.—Co-added images taken at rotator offsets of 0 and -30° . Structure in the PSF is very similar between the two, and rotates along with the rotator offset. This indicates that the aberrations leading to the structure are in the telescope optics. For this reason, "rolling" the telescope by using the derotator will not help remove the PSF structure.

To quantify the achievable contrast limit in an image, a boxcar smoothing of the image over a 3×3 pixel box was applied, and annuli around the central pixel were analyzed to estimate the variation in flux in an aperture approximately the size of the PSF. For each annulus the standard deviation of all the included pixels was calculated. A threshold of 5 times the standard deviation is a reasonable estimate of the limiting contrast versus separation. A plot of these values is shown in Figure 4 for three separate images. The solid line shows the limit for a raw PSF, the dashed line shows the limit for an unsharp-masked (5×5 pixel boxcar smoothed) image, and the dot-dashed line shows the result for a PSF-subtracted image, which has also been unsharp masked.

Figure 4 verifies that PSF subtraction using sky rotation will be a useful approach for high dynamic range imaging in the M band. We expect to achieve contrast limits of roughly 11 mag at 1" and 13 mag by 1".5. While sky rotation did not allow us to take sufficient data to use this method for Vega, these initial data



5 sigma Contrast Limit

FIG. 4.—Contrast limit vs. separation. The 5 σ contrast limit in *M*-band magnitudes are plotted for Clio on the MMT for a PSF image (*solid line*), an unsharp masked image (*dashed line*), and a PSF-subtracted image (*dash-dotted line*).

suggest that observing an object through significant sky rotation will be the most straightforward way of maximizing the detection sensitivity of a faint companion.

5. DISCUSSION

Companions to Vega have been searched for using several other systems, including the Palomar AO system (Metchev et al. 2003), the Keck AO system (Macintosh et al. 2003a, 2003b), and the Gemini AO system (Marois et al. 2006, hereafter M06). At 7", the distance of a likely perturbing companion for the cold dust, Metchev et al. estimate a 5 σ limit of H = 14.2, which corresponds to a 19 M_J planet, using the models of Burrows et al. (1997). For Macintosh et al. the limiting magnitude at the *K* band is K = 17.3 or $10 M_J$. M06 made significant improvement on the *H*-band limit of Metchev et al. through the use of angular differential imaging (ADI). They present a 12,000 s observation that achieves a 5 σ limit of H = 19.5 at 7", an impressive limit which is equivalent to a 4 M_J planet.

Analysis of the Clio data, as shown in Figures 1 and 2, suggests that at 7" separation we are constrained by background noise to a limit of about 0.8 mJy (13.3 mag), or 7 M_J for an approximately 11 minute observation. This limit is similar for separations as close as approximately 2".5 from the star. The background noise appears random in this region, allowing us to extrapolate the expected sensitivity for longer integrations. For example, a similarly long exposure of 12,000 s at the *M* band can reasonably be expected to achieve a limiting magnitude of 0.18 mJy (14.8 mag), corresponding to a companion limit of approximately 2.5 M_J around Vega, and which would be detectable as close as a few arcseconds from the star.

5.1. Implications for Planet Detection

The relative contrast versus separation achievable in the M band presented in Figure 4 is similar to what has been achieved with either *Hubble Space Telescope* (*HST*) or the best ground-based adaptive optics systems in the near-infrared. For example, M06 demonstrate a limit of 11.1–11.9 mag at 0.8 separation on Gemini using the Altair system at the H band. Using spectral difference imaging on the VLT, a 5 σ limit of 11.0 has been achieved (Biller et al. 2005). Figure 4 shows a 5 σ limit of 11 at 1". Scaled

Distance (pc)	Band	Age (Gyr)	Star Magnitude	Planet Magnitude	Contrast	Minimum Separation (arcsec)	Minimum Separation (AU)
5	М	5	2	14.2	12.2	1.2	6
	Н	5	2	22.4	20	8	40
20	M	0.5	5	14.9	9.9	0.8	16
	Н	0.5	5	19.2	14.2	1.5	30
50	M	0.1	7	15.5	8.5		
	Н	0.1	7	18.4	11.4	0.8	81

TABLE 2 Expected Minimum Detection Angle for a 10 $M_{\rm J}$ Planet around a Solar Twin at H and M Bands

Notes.—The conversion from contrast to minimum separation for the M band is from Fig. 4. The conversion for the H band is from M06.

to the same telescope diameter, these three observational results have similar limiting contrast levels.

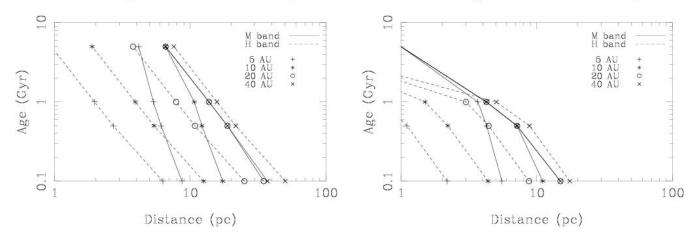
In addition to the improved PSF stability, the relative contrast of Jupiter-like companions at longer wavelengths provides a significant advantage for the approach outlined here. Burrows et al. (2003) predict that the contrast for a 10 M_J planet around Vega is 15.9 mag at H band, while it is only 12.7 at M band, a difference of 3.2 mag. For a 3 M_J planet, the H- and M-band magnitudes are predicted to be 20.2 and 14.7, respectively. For close separations where PSF subtraction dominates the achievable contrast, observing at longer wavelengths can lead to detection of significantly less massive companions. To be clear, the current data presented here do not yet demonstrate detectability at this level, but the PSF subtractions carried out show that the contrasts are reachable in tandem with observations of an object spanning a significant change in parallactic angle.

The observations presented here illustrate that an optimized imager, when used with an integrated adaptive optics system, can provide detection capabilities similar to, or even better than, near-infrared techniques. The useful detection for the *M* band is outside of $\sim 3 \lambda/D$, or 0".5 for the MMT, where a contrast of approximately 9 mag is achieved. Inside this region the shorter wavelengths could likely be used to detect relatively bright companions ($\Delta m < 5$) that would be more difficult to extract at the lower spatial resolution achieved at the *M* band. Outside of this region, the observations presented here show that $3-5 \mu m$ imaging can provide contrast ratios similar to those from separation.

Coupled with the more favorable contrast in the spectral window, this makes the search for planetary mass companions at longer wavelengths an attractive alternative to near-infrared observations.

Although the L' and M bands provide interesting spectral regions to find planets, optimum candidate stars for observations at the near-infrared versus L' and M will likely be quite different. For the *H* band, typical surveys have focused on very young stars (<300 Myr) in order to be able to detect the planets while they are still relatively hot and thus bright in the H band. With an L'and M-band survey the most attractive stars become the most nearby stars, although youth is still important. Several examples of this are detailed in Table 2 for typical targets in the H and M bands. A 10 $M_{\rm J}$ planet in a wide orbit around a 5 Gyr old solar twin at 5 pc ($M_s \simeq H_s \simeq 2$) would have $M_p = 14.2$ (Baraffe et al. 2003, hereafter Ba03). The contrast of 12.2 mag would allow detection, from Figure 4, outside of 1"2 or 6 AU. The H-band flux from Ba03 would be $H_p = 22.4$, a contrast of 20 mag, which is detectable outside of 8" or 40 AU according to the ADI contrast limit of M06. For a 0.5 Gyr old G star at 20 pc ($M_s \simeq H_s \simeq 5$) Ba03 predict a 10 M_J planet to be $M_p = 14.9$, a contrast of 9.9 mag, which would be detectable outside of 0.78, or 16 AU. The *H*-band flux for such a planet is $H_p = 19.2$, a contrast of 14.2 mag, which is detectable outside of 1".5 or 30 AU according to the ADI limit of M06. If we consider an even younger solar twin at 50 pc and an age of 0.1 Gyr, a 10 $M_{\rm J}$ planet, at $M_p = 15.5$, would not be detectable at the *M* band, and for the H band be detectable outside 0".8 or 81 AU. These examples

Distance-Age Limits for a 3 M_J planet



Distance-Age Limits for a 10 M_J planet

Fig. 5.—Distance-age limits for 10 and 3 M_J planets around a solar twin. The limits are shown for orbits of 5, 10, 20, and 40 AU. The contrast vs. separation for the M band are taken from Fig. 4, with the exception that a sky background limit of M = 14.8 is used for larger separations. The contrast vs. separation for the H band is from M06, with a sky background limit of H = 23.5 assumed.

show the general trend: *H*-band detections on large-aperture telescopes are best suited to young stars (which are necessarily further away), while $3-5 \mu m$ detection is likely to be more sensitive for nearby stars.

The parameter space of detection for the M versus the H band can be visualized by plotting the limiting distance and age at which a planet can be detected for various orbits around a solar twin. Figure 5 shows two plots that illustrate this for a 10 M_J and a 3 M_J planet. The contrast limits are taken from Figure 4 for the M band and from M06 for the H band. A sky background limit of M = 14.8 and H = 23.5 is also assumed. Planet fluxes in each band are taken from Ba03. The M band is preferable to the H band on the plots unless the planet is in a wide (40 AU) orbit or is a young system (<0.1 Gyr).

Planet detection in the L'- and M-atmospheric windows requires a large-aperture telescope and adaptive optics system, as do techniques at shorter wavelengths. However, for L'- and M-band observations, a noise floor at larger separations arises from the sky background. For the parameters of our observations of Vega, this arises at $\sim 10 \lambda/D$. Thus, for L'- and M-band observations, a large aperture benefits in not only the sharper image quality (and thus a closer inner working distance) but also fainter limiting sensitivity for companions at larger separations. This makes AO imaging at the L' and M bands increasingly favorable for larger aperture telescopes. For example, the Large Binocular Telescope with its 2×8.4 m aperture in coherent imaging mode will have 3.4 times the collecting area and concentrate the light into a PSF that has an angular area 3.7 times less than that of the MMT PSF. This translates to a limiting magnitude improvement of 2.7. Extrapolating from the MMT sensitivity, we could expect to be sensitive to a 2 $M_{\rm I}$ planet around a 1 Gyr old star at 10 pc.

Next-generation telescopes, such as the Thirty Meter Telescope³ or the Giant Magellan Telescope,⁴ can improve these limits further through closer inner working distances and fainter sensitivities, but it will be important for these telescopes to have adaptive optics that are integrated into the telescope, such as the deformable secondary approach being planned for the GMT.

6. CONCLUSION

We present initial *M*-band observations of Vega with a newly commissioned thermal infrared camera, Clio, on the 6.5 MMT, using the integrated deformable secondary mirror. The observation constrains a planet to be less than 7 M_J in the orientation expected from the dust structure to approximately 2".5 (20 AU) from the star. Extrapolation to longer integration times suggests that observations similar to these will be sensitive to wide companion planets at the top end of the mass distribution seen by radial velocity searches for nearby stars.

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- ³ See http://www.tmt.org.
- ⁴ See http://www.gmto.org.

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